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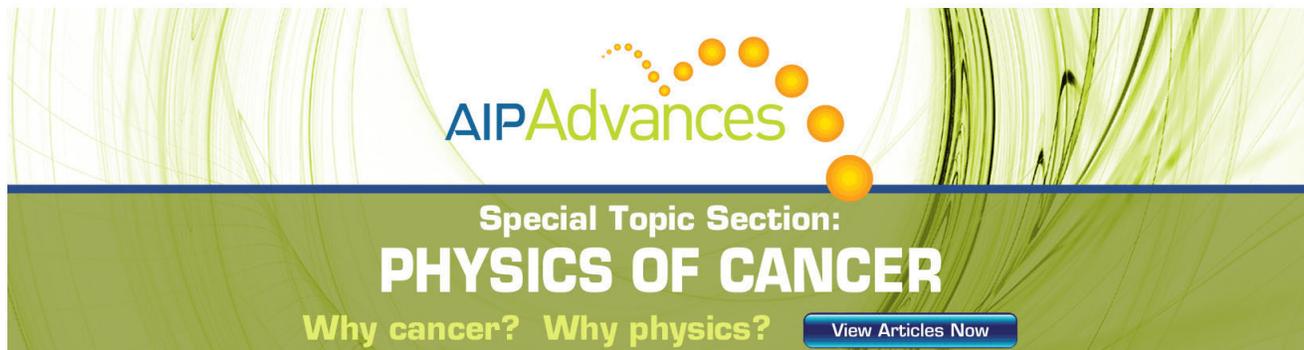
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Enhanced optical transmission through tapered metallic gratings

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To achieve non-resonant behavior, we propose a metallic grating device with linearly tapered slits. The tapering provides a gradual impedance variation from the entrance to the exit of the slits, leading to broadband and wide-angle enhanced transmission in the infrared. In addition, the light is strongly localized and enhanced at the slit exits, in contrast with straight slits. We describe the phenomenon with a transmission line model, which is in accordance with rigorous simulations. © 2012 American Institute of Physics. [<http://dx.doi.org/10.1063/1.4729005>]

The extraordinary optical transmission (EOT) phenomenon was observed in an opaque metallic film with a periodic array of subwavelength holes in 1998.¹ Since then tremendous investigation has been dedicated to its mechanism,²⁻⁵ elucidating the role of surface plasmon dispersion⁶ and Fabry-Perot type cavity modes.⁷ EOT offers strongly localized field enhancements in addition to a large transmission, with potential applications in nonlinear optics, sensing etc. However, due to the resonant mechanisms involved often a relatively narrow bandwidth is affected, which is unsuitable for a range of applications. On the other hand, broadband transparency has been achieved in related structures, e.g., using oblique incidence TM polarization,^{5,8} or by connecting rectangular apertures with narrower slits.⁹ In the latter example, the thin gratings operate at wavelengths beyond the Fabry-Perot resonances of guided modes in the slits. This regime of non-resonant devices has only recently attracted attention, but it is indispensable e.g., for broadband light harvesting devices.

Therefore, we propose metallic gratings with tapered slits, which offer a much broader transmission window for TM polarized light. By gradually varying the impedance from input to output plane, we effectively destroy the Fabry-Perot type resonant conditions of guided modes in the slit, yielding a non-resonant and thus broadband and wide-angle large transmission in the infrared. In addition, the localization of the field is confined to one plane of the structure, instead of over the whole lossy waveguide as in the traditional structures with straight sidewalls.

Analysis is performed with rigorous calculations (finite element method (FEM)¹⁰) and semi-analytic methods (transmission line (TL) theory⁸). Other more complex models and theories were proposed for metallic aperture devices,^{11,12} but under certain conditions the efficient TL model provides effective characterization. To illustrate the effects of tapering we assume surrounding, substrate, and slit material to be air. The standard geometry and the tapered grating are shown in Figs. 1(a) and 1(b), respectively. The grating with tapered slits is characterized by the thickness t , period p , and widths

w_{in} and w_{out} at the entrance and exit of the slit, respectively. The metal grating is assumed to be silver.¹⁵ Fig. 2 shows the calculated TM transmission T with normal incidence ($\theta = 0$) for grating structures with different sizes, with both FEM and TL calculations. For various parameters, the tapered grating leads to a dramatic enhancement of the transmission over a broadband wavelength range, comparing the straight sidewall case (Fig. 2, blue curves) with the tapered cases. In addition, we notice that the tapered transmission seems to average the Fabry-Perot resonant peaks for the non-tapered geometry, leading to non-resonant transmission. This averaging effect is clear for the smaller wavelengths, and becomes more pronounced for larger w_{in} (more “open” gratings).

We introduce now the employed TL approach, which was developed for non-tapered slits.⁸ Surrounding and substrate are modeled as semi-infinite TLs characterized by a wavenumber (β_{in} , β_s , respectively) and a characteristic impedance per unit length (Z_{in} , Z_s , respectively), see circuit in Fig. 1(a). The wavenumber is given by $\beta_u = k_0 n_u \cos \theta$, the characteristic impedance is $Z_u = Z_0 p \cos \theta / n_u$, $u = in$ (surrounding), s (substrate), and $Z_0 = \sqrt{\mu_0 / \epsilon_0}$ is the vacuum impedance. The grating layer is treated as a TL characterized by β and Z with finite length t . Given that the metal width d ($= p - w$) is large enough to avoid coupling between slits, the wavenumber β is determined by the dispersion relation of a metal-insulator-metal waveguide:¹³

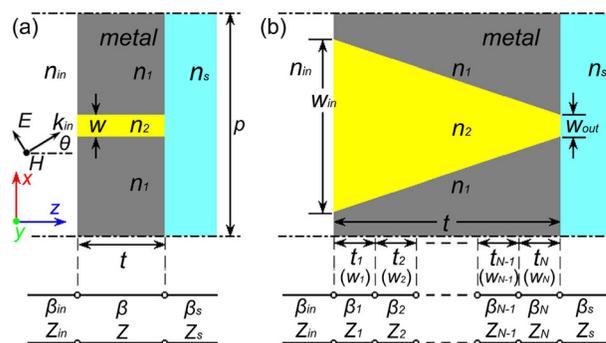


FIG. 1. Schematics of 1D metallic grating with (a) straight slits, and (b) linearly tapered slits, respectively, together with illustrations of the TL model underneath. The gratings are periodic in the x -direction.

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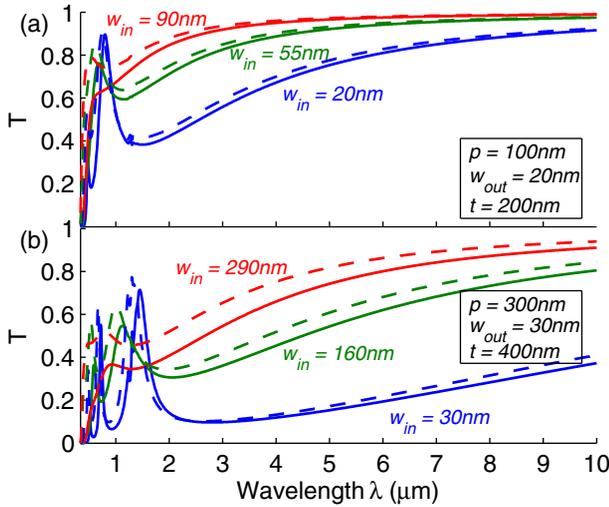


FIG. 2. Transmission T for gratings with different sizes (i.e., different slit width w_{in} and w_{out} , grating period p , slit thickness t). The solid lines are from numerical simulation by FEM. The dashed lines are from TL model.

$$\tanh\left(\sqrt{\beta^2 - k_0^2 n_2^2} w/2\right) = -\frac{\sqrt{\beta^2 - k_0^2 n_1^2} n_2^2}{\sqrt{\beta^2 - k_0^2 n_2^2} n_1^2}. \quad (1)$$

Then, the characteristic impedance per unit length can be expressed as⁸

$$Z = w\beta/(\omega\epsilon_0). \quad (2)$$

The tapered grating is modeled as a series of cascaded TLs, with a staircase multilayer characterized by local widths w_i and thicknesses t_i (circuit in Fig. 1(b)). To calculate the transmission T of the cascaded network one can use $ABCD$ parameters (also known as transfer matrix).¹⁴ The $ABCD$ parameters (M) of the whole network are a multiplication of each TL in the network: $M = [A_{tot} B_{tot}; C_{tot} D_{tot}] = \prod M_i = \prod [A_i B_i; C_i D_i]$, where $A_i = D_i = \cos(\beta_i t_i)$, $B_i = jZ_i \sin(\beta_i t_i)$, $C_i = jY_i \sin(\beta_i t_i)$, $Y_i = 1/Z_i$, $i = 1, 2, \dots, N$, ($i = 1, 2, \dots, N$). Therefore T is expressed by

$$T = \frac{4|Y_{in}|^2 \text{Re}(Y_s)}{|m_1 Y_{in} + m_2|^2 \text{Re}(Y_{in})}, \quad (3)$$

where $m_1 = A_{tot} + B_{tot} Y_s$, $m_2 = C_{tot} + D_{tot} Y_s$, $Y_u = 1/Z_u$, $u = in, s$. Upon derivation of Eq. (2), uniform electric fields in the slit and negligible coupling between the slits are assumed,⁸ so the TL model works well with small slit widths (w) and large enough metal widths (d). In addition, diffraction in the surrounding and substrate, and surface modes at the grating planes are not considered in the TL model, so the period of the grating cannot be larger than the wavelength.

The plots in Fig. 2 show a good agreement between the rigorous full-wave FEM simulations and the TL model. Here, in the TL model 10 layers are used and already provide a good convergence. More in detail, agreement is better for small periods (Fig. 2(a)), and for straight slits (blue curves in Figs. 2(a) and 2(b)). The deviation is mainly caused by plasmon modes at the slit corners, which are not taken into account in the TL model. Deviation increases when w_{in}

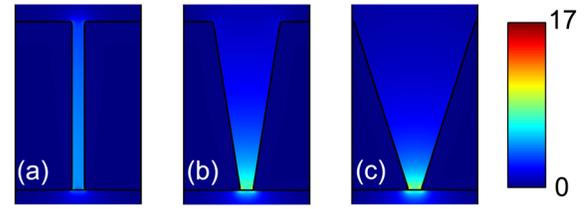


FIG. 3. E field amplitude (normalized by the incident field amplitude E_0) distribution at $\lambda = 5 \mu\text{m}$ for gratings ($t = 400 \text{ nm}$, $w_{out} = 30 \text{ nm}$, $p = 300 \text{ nm}$) with (a) $w_{in} = 30 \text{ nm}$, (b) $w_{in} = 160 \text{ nm}$, and (c) $w_{in} = 290 \text{ nm}$.

increases, since the coupling between slits becomes stronger, and uniformity of the E field in the slits diminishes. However, the TL model still predicts the trend of T very well, except for an overestimation at smaller wavelengths.

Based on the large transmission enhancement by tapering, we can imagine that the light is slowly squeezed from the entrance of the taper on to the narrower exit slit. Therefore, the field is expected to be gradually enhanced. To confirm, we plot the electric field spatial distribution (Fig. 3) at wavelength $5 \mu\text{m}$ for gratings corresponding to Fig. 2(b). For other grating sizes, the field spatial distribution profiles are very similar. We notice indeed that the E amplitude gradually increases towards the exit of the slit by tapering, contrast (a) to (b) and (c). In addition, the fields reach their maximum exactly at the exit of the tapered slit, (c) and the maximum value is much larger than in the straight sidewall case.

To examine the field enhancement properties, Fig. 4 shows the average normalized electric field as a function of wavelength at the slit exit (same gratings as in Fig. 3). Broadband field enhancement is obtained, and it increases as the taper becomes wider (increasing w_{in}). Meanwhile, the spectrum of the normalized field is similar to the corresponding transmission spectrum (Fig. 2(b)). Therefore, the near-field at the exit of the slit and the transmission have a similar spectrum, with deviations from evanescent wave components. In addition, for larger wavelengths, we notice that the normalized electric field arrives quite close to the dashed-dotted line, which is the ratio of p/w_{out} , similar to Ref. 9. The average field decreases towards smaller wavelengths, since there is less light transmitted. However, by tapering it is still possible to obtain a local field enhancement far beyond the ratio of p/w_{out} , instigated by the sharp corner at the exit of the slit, as the field maxima shown in Fig. 3. Therefore, tapering offers a strong control over the field

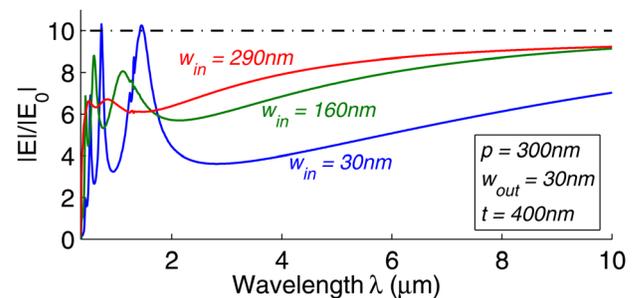


FIG. 4. Average E field amplitude (normalized by the incident field amplitude E_0) versus wavelength at the exit of slit for gratings ($t = 400 \text{ nm}$, $w_{out} = 30 \text{ nm}$, $p = 300 \text{ nm}$) with different widths: $w_{in} = 30 \text{ nm}$ (blue), $w_{in} = 160 \text{ nm}$ (green), and $w_{in} = 190 \text{ nm}$ (red).

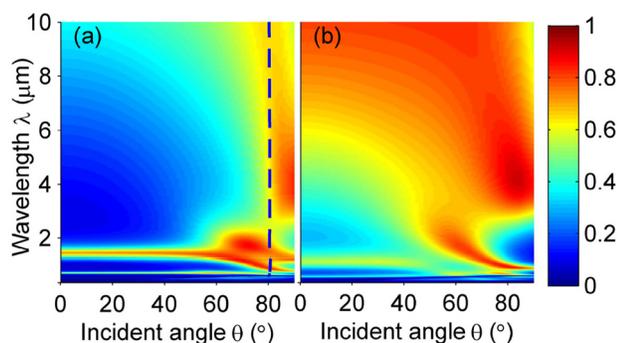


FIG. 5. Angular response for gratings ($t = 400$ nm, $p = 300$ nm, $w_{out} = 30$ nm) with (a) $w_{in} = 30$ nm (straight sidewalls), (b) $w_{in} = 160$ nm (tapered sidewalls). The dashed line in (a) is the plasmonic Brewster angle.

enhancement profile, by tailoring the transmission spectrum, the value of p/w_{out} and the corner sharpness.

Finally, we compare the angular response of the straight sidewall gratings (Fig. 5(a)) with the tapered gratings (Fig. 5(b)). For the straight sidewall (Fig. 5(a)), we observe that the transmission gradually increases with the angle, and reaches its maximum near the dashed line around 80° , which is the Brewster angle θ_B determined by impedance matching between the surrounding and the straight slit.⁸ For the tapered case (Fig. 5(b)), we see that the overall transmission is significantly larger. In addition, we still observe Brewster angle behavior, i.e., the transmission increases with incident angle, although the impedance in the slit is not uniform.

In experimental realizations, a different substrate may be necessary. Our calculations show the validity of the tapering idea for $n_s \neq n_i$, e.g., with a glass substrate, except that the transmission at larger wavelength is determined by that of light through a semi-infinite surrounding and substrate ($T = 4n_s n_{in} / (n_s + n_{in})^2$).

In summary, we demonstrate the concept of tapering in metallic gratings to achieve broadband and wide-angle transmission. In addition, the taper provides a strong enhancement and localization of light at the exit of the slit, useful for

applications such as nonlinear optics, light harvesting, sensing,¹⁶ and emission enhancement. The TL model gives a very efficient characterization of the tapering effect, therefore it can be used to assist more complex designs, e.g., with parabolic shapes. Finally, the transmission at smaller wavelengths is further improved by tilting the incident light, giving rise to a plasmonic Brewster type effect. This type of structure therefore expands the field of non-resonant gratings and broadband transmission.

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¹T. W. Ebbesen, H. J. Lezec, H. F. Ghaemi, T. Thio, and P. A. Wolff, *Nature (London)* **391**, 667 (1998).

²H. Ghaemi, T. Thio, D. Grupp, T. Ebbesen, and H. Lezec, *Phys. Rev. B* **58**, 6779 (1998).

³J. Porto, F. García-Vidal, and J. Pendry, *Phys. Rev. Lett.* **83**, 2845 (1999).

⁴L. Martín-Moreno, F. García-Vidal, H. Lezec, K. Pellerin, T. Thio, J. Pendry, and T. Ebbesen, *Phys. Rev. Lett.* **86**, 1114 (2001).

⁵X.-R. Huang, R.-W. Peng, and R.-H. Fan, *Phys. Rev. Lett.* **105**, 243901 (2010).

⁶W. L. Barnes, A. Dereux, and T. W. Ebbesen, *Nature (London)* **424**, 824 (2003).

⁷Q. Cao and P. Lalanne, *Phys. Rev. Lett.* **88**, 057403 (2002).

⁸A. Alù, G. D’Aguanno, N. Mattiucci, and M. Bloemer, *Phys. Rev. Lett.* **106**, 123902 (2011).

⁹G. Subramania, S. Foteinopoulou, and I. Brener, *Phys. Rev. Lett.* **107**, 163902 (2011).

¹⁰COMSOL Multiphysics, a commercial FEM software. Available at <http://www.comsol.com>.

¹¹F. Medina, F. Mesa, and R. Marqués, *IEEE Trans. Microwave Theory Tech.* **56**, 3108 (2008).

¹²V. Delgado and R. Marqu, *Opt. Express* **19**, 25290 (2011).

¹³E. N. Economou, *Phys. Rev.* **182**, 539 (1969).

¹⁴D. M. Pozar, *Microwave Engineering* (Wiley, Danvers, 2005).

¹⁵E. D. Palik, *Handbook of Optical Constants of Solids* (Academic, New York, 1985).

¹⁶M. Schnell, P. Alonso-González, L. Arzubiaga, F. Casanova, L. E. Hueso, A. Chuvilín, and R. Hillenbrand, *Nat. Photonics* **5**, 283 (2011).